

Neutrino mass generation in asymptotically safe gravity

Gustavo P. de Brito

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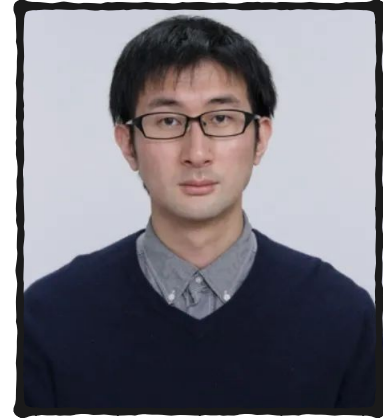
In collaboration with:



Astrid Eichhorn



Antônio D. Pereira



Masatoshi Yamada

To appear: 2504.XXXXX

Interplay between gravity and matter in asymptotic safety

Towards experimental tests of quantum gravity

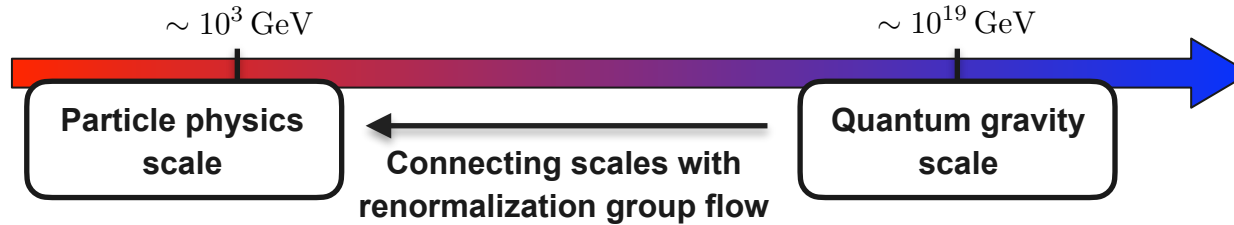
➔ How to connect quantum gravity with experimental tests?



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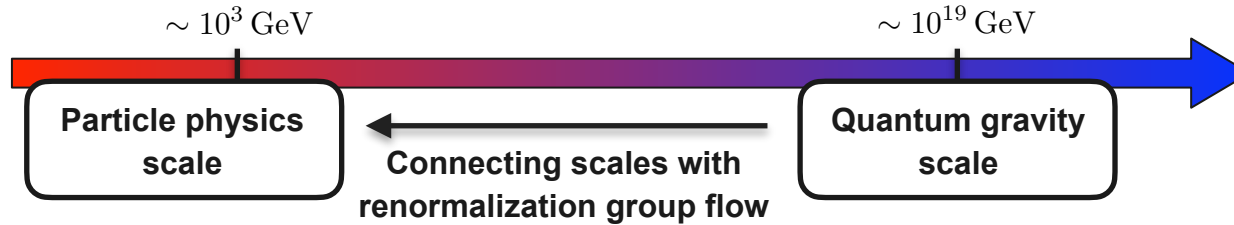
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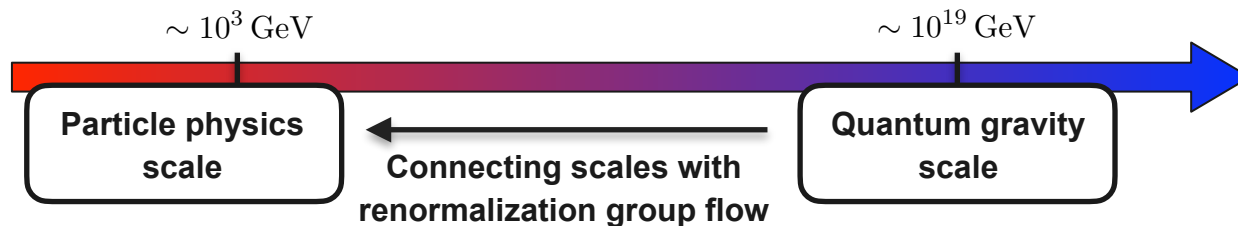
Exciting results in the past 15 years

- ▶ **Higgs mass from ASQG**
Shaposhnikov, Wetterich (2009), ...
- ▶ **Solution to the hypercharge triviality problem** Christiansen, Eichhorn (2017), ...
- ▶ **UV completion of SM**
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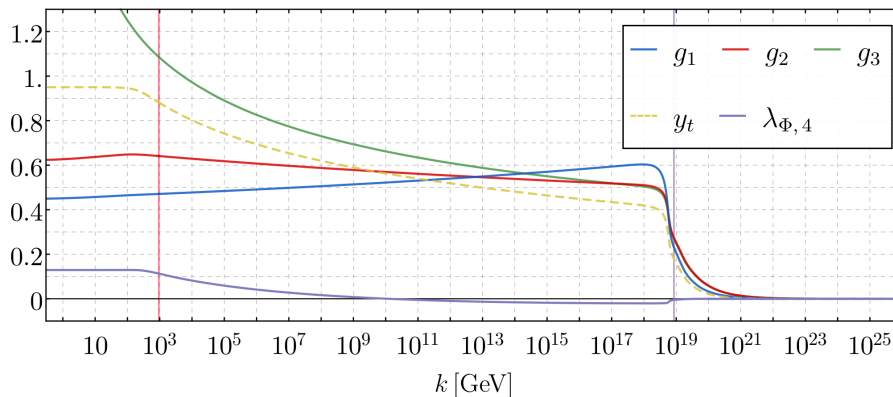
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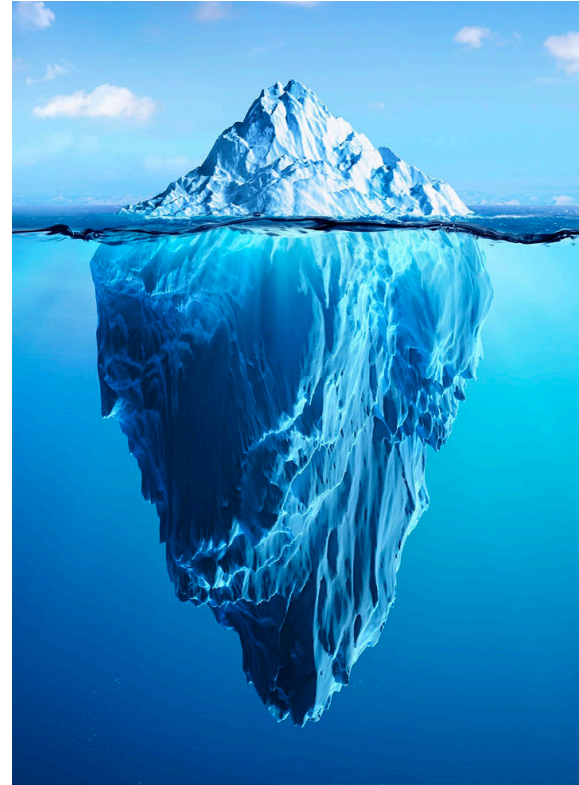
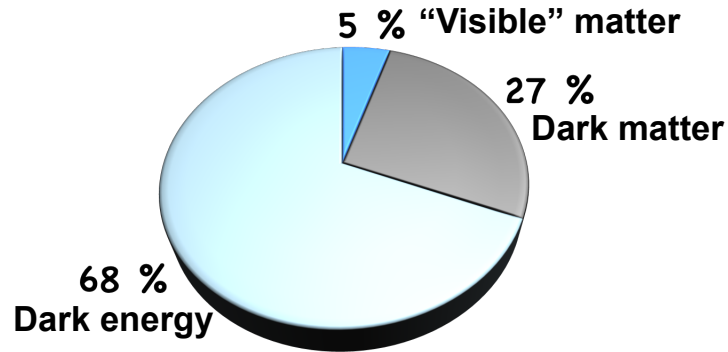
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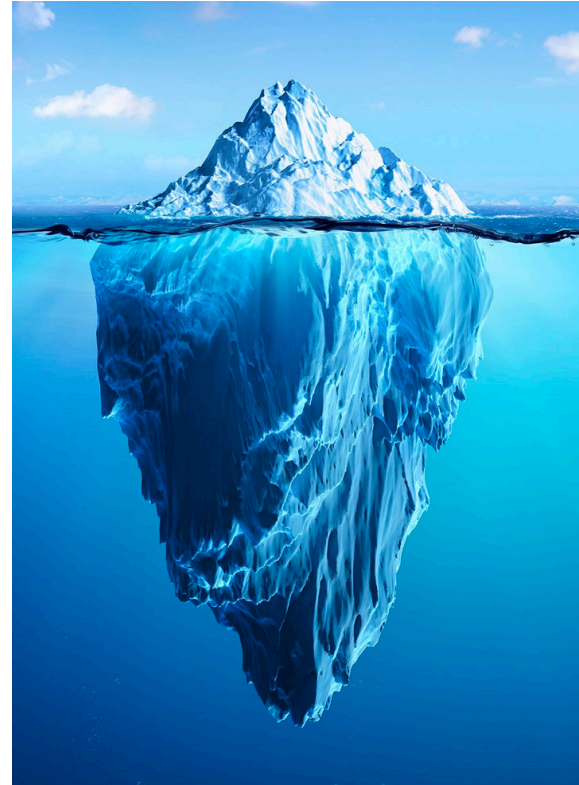
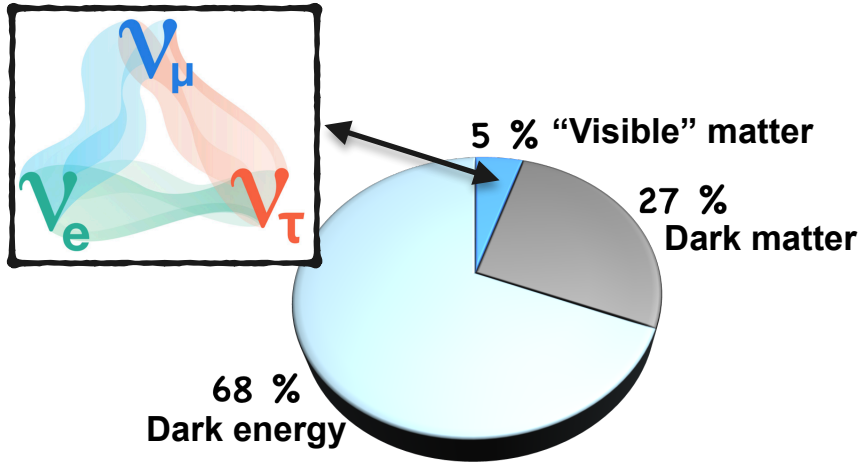
Missing pieces of the Standard Model

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Quick detour in the dark side of asymptotically safe gravity

- Asymptotically safe gravity can lead to theoretical constraints on Dark Matter candidates
- Active searches of dark matter could work as indirect tests of asymptotically safe gravity



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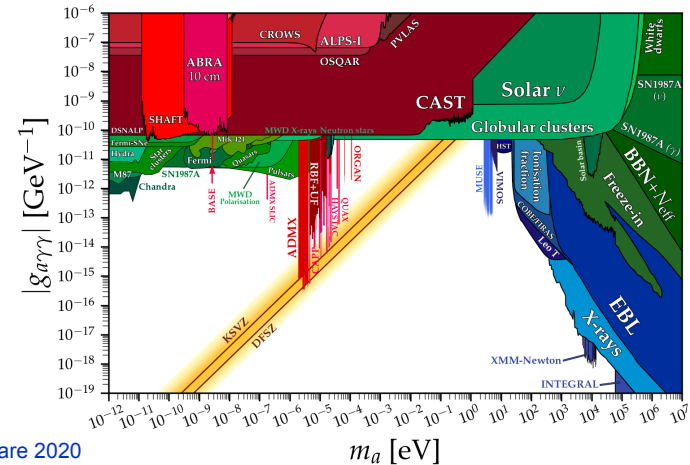
Axion-Like Particles (ALPs) in ASQG?

- Typical ALP-photon interaction:

$$\mathcal{L}_{\gamma\text{-ALP}} = \frac{g_{a\gamma\gamma}}{4} \varphi F_{\mu\nu} \tilde{F}^{\mu\nu}$$

- Popular candidate for ultra-light (non-thermally produced) dark matter

Reviews on ALPs: Arias et.al.(2012), Ringwald (2012,2014), Irastorza, Redondo (2018), Irastorza (2021), ...



O'Hare 2020
(<https://cajohare.github.io/AxionLimits/>)

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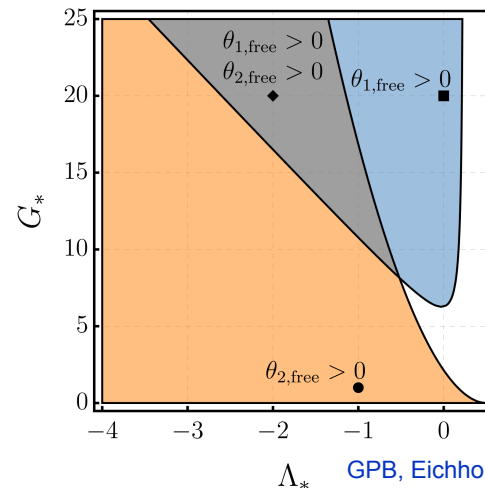
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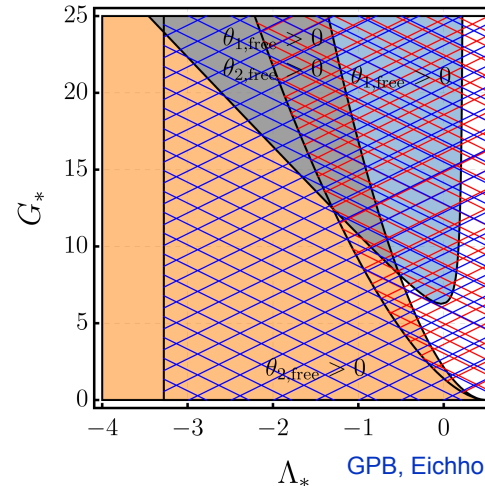
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Ruling out vector dark matter models in ASQG

- DM via hidden gauge sector

$$\mathcal{L}_{\text{DM}}^{\text{UD}(1)} = -\frac{1}{4}V_{\mu\nu}V^{\mu\nu} + (D_\mu S)^*(D^\mu S) - V(\Phi, S)$$

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- Viable phenomenology

[Frandsen, et.al. \(2023\)](#)

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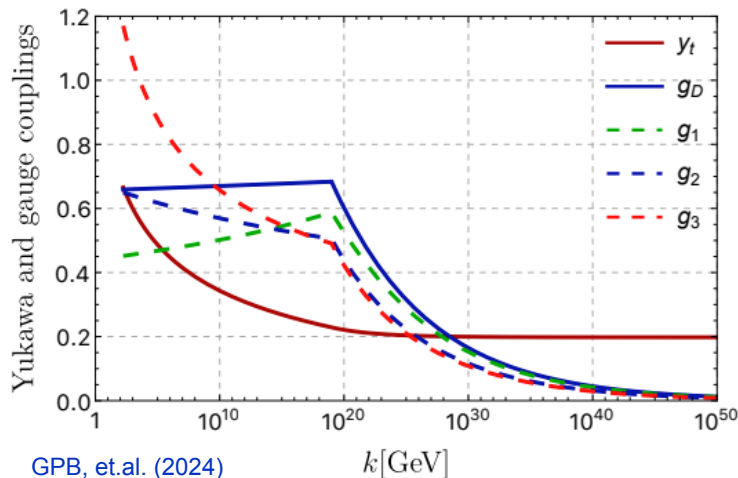
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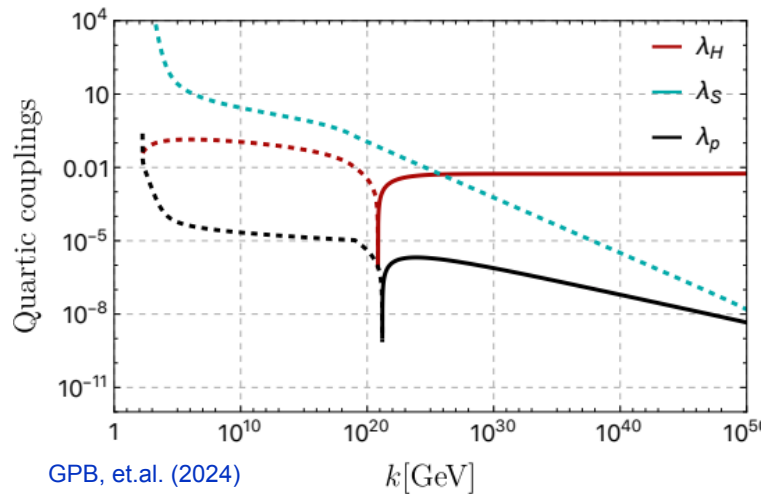
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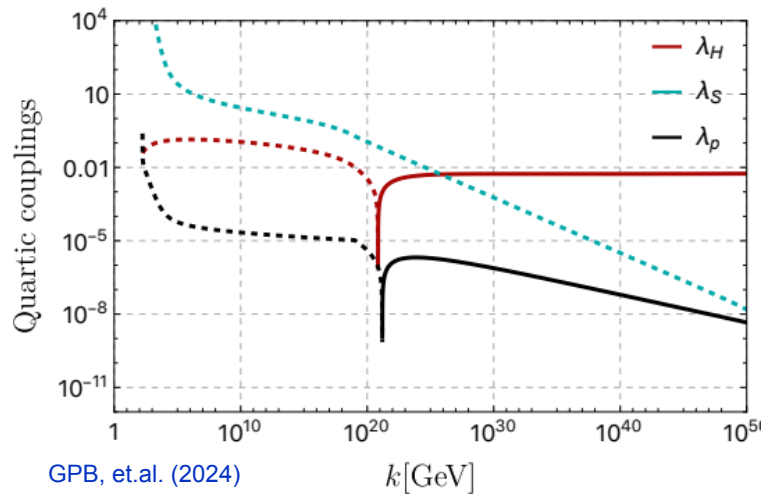
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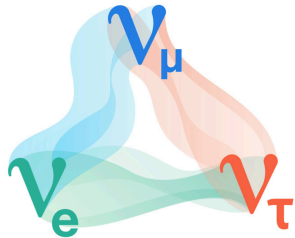


Massive Neutrinos: another missing ingredient in the SM

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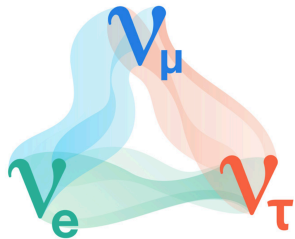


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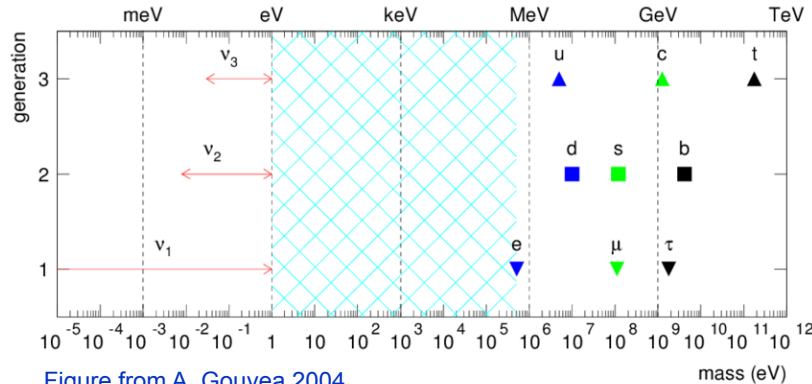
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→ Explanation: Neutrinos are massive particles (with mass basis different from flavour basis)

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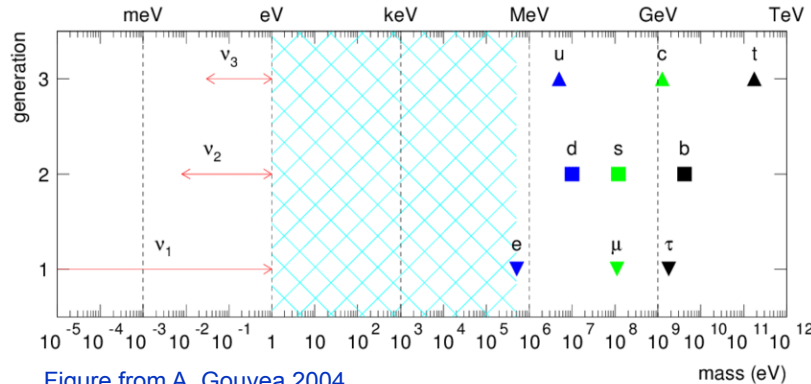


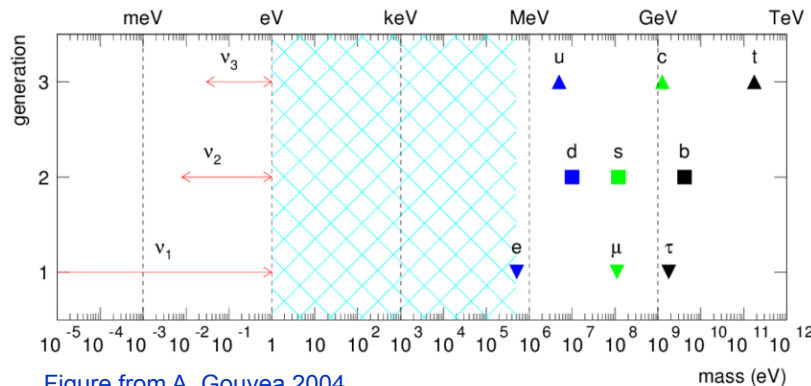
Figure from A. Gouvea 2004

* Recent upper bounds from DESI survey

$$\sum_{\nu} m_{\nu} < 0.06 \text{ eV (inverted hierarchy)}$$
$$\sum_{\nu} m_{\nu} < 0.10 \text{ eV (normal hierarchy)}$$

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- What is the origin of neutrino masses?
- Is it possible to explain the the large hierarchy between neutrino masses and other fermion masses?

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Neutrino Masses in Asymptotically Safe Gravity

- Many alternatives in the market (Dirac neutrinos, See-Saw, Weinberg operator...). Is there a theoretical principle to select a few of them?
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Naturally small Yukawa couplings from trans-Planckian asymptotic safety

Kamila Kowalska, Soumita Pramanick and Enrico Maria Sessolo

ABSTRACT: In gauge-Yukawa systems embedded in the framework of trans-Planckian asymptotic safety we discuss the dynamical generation of arbitrarily small Yukawa couplings driven by the presence of a non-interactive infrared-attractive fixed point in the renormalization group flow. Additional ultraviolet-attractive fixed points guarantee that the theory remains well defined up to an infinitely high scale. We apply this mechanism to the Yukawa couplings of the Standard Model extended with right-handed neutrinos, finding that asymptotically safe solutions in agreement with the current experimental determination of the masses and mixing angles exist for Dirac neutrinos with normal mass ordering. We generalize the discussion by applying the same mechanism to a new-physics model with sterile-neutrino dark matter, where we generate naturally the feeble Yukawa interaction required to reproduce via freeze-in the correct relic abundance.

Dynamically vanishing Dirac neutrino mass from quantum scale symmetry

Astrid Eichhorn^a, Aaron Held^{b,c, *}

A B S T R A C T

We present a mechanism which drives Dirac neutrino masses to tiny values along the Renormalization Group flow, starting from an asymptotically safe ultraviolet completion of the third generation of the Standard Model including quantum gravity. At the same time, the mechanism produces a mass-splitting between the neutrino and the quark sector and also generates the mass splitting between top and bottom quark. The mechanism hinges on the hypercharges of the fermions and produces a tiny neutrino Yukawa coupling, because the right-handed neutrino is sterile and does not carry hypercharge.

See also: A. Held Ph.D. thesis, 2019

See also:

Domenech, Goodsell and Wetterich, 2021
Chikaballi, Kowalska and Sessolo, 2023

For massive neutrinos with see-saw scale

Neutrino Masses in Asymptotically Safe Gravity

In the rest of this talk, I will discuss three scenarios for massive neutrinos in ASQG

- **Scenario I:**
Massive neutrinos from Weinberg operator
(without right-handed neutrinos)
- **Scenario II:**
Majorana massive neutrinos from see-saw mechanism (type I)
- **Scenario III:**
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- ➔ Compatible with asymptotically safe gravity
- ➔ No upper or lower bound from this scenario
- ➔ We can tune the relevant parameters to embed this scenario into asymptotically safe gravity

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Working hypothesis:

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Strong evidence from pure gravity and gravity + “minimal matter”

Falls, Litim, Nikolakopoulos and Rahmede, 2013

Falls, Litim and Schröder, 2018

Eichhorn, Lippoldt, Pawłowski, Reichert and Schiffer, 2018

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- Usually considered as part of the SMEFT (as a lepton number violating term)

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Can we make sense of this scenario beyond the EFT paradigm?

(Gravity induced UV completion?)

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➔ We look at the beta function of ζ including quantum gravity contributions

$$\beta_\zeta = \zeta \left(1 - \frac{3}{16\pi^2} g_2^2 + \frac{3}{8\pi^2} (y_t^2 + y_b^2 - y_\tau^2/6) + \frac{1}{4\pi^2} \lambda_H + \frac{17}{18\pi} G \right)$$

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* The inclusion of a cosmological constant does not change this picture, unless we allow “unrealistic” fixed point values of the gravitation couplings (e.g. $G^* > 30$)

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- ➔ The Weinberg operator scenario does not become UV complete under the impact of gravity
- ➔ Asymptotically safe gravity seems to require new degrees of freedom in the neutrino sector
- ➔ This results does not rule exclude the Weinberg operator in a EFT setting

Majorana Neutrinos and See-Saw Mechanism

- A popular scenario for massive neutrinos is based on the see-saw mechanism

$$\mathcal{L}_\nu \supset \frac{m_R}{2} (\bar{\nu}_R \nu_R^C + \text{h.c.}) + y_\nu \left(\bar{L} \tilde{H} \nu_R + \text{h.c.} \right) \xrightarrow{\text{SSB}} \mathcal{L}_\nu \supset \frac{m_R}{2} (\bar{\nu}_R \nu_R^C + \text{h.c.}) + m_D (\bar{\nu}_L \nu_R + \text{h.c.}) + \dots$$
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$$\rightarrow \text{Mass matrix } M_\nu = \begin{pmatrix} 0 & m_R \\ m_R & m_D \end{pmatrix} \quad m_D = \frac{1}{\sqrt{2}} y_\nu v_H$$

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→ Mass matrix $M_\nu = \begin{pmatrix} 0 & m_R \\ m_R & m_D \end{pmatrix}$ $m_D = \frac{1}{\sqrt{2}} y_\nu v_H$

- The eigenstates of the mass matrix are Majorana fermions (neutrinos ~ anti-neutrinos)

→ Smoking gun signature for Majorana neutrinos: neutrinoless double beta decay

* Searches at various experiments:
NEMO-3;
NEXT-100; KamLAND-Zen;
EXO-200; CUORE; GERDA

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$$\Rightarrow \text{Mass matrix } M_\nu = \begin{pmatrix} 0 & m_R \\ m_R & m_D \end{pmatrix} \quad m_D = \frac{1}{\sqrt{2}} y_\nu v_H$$

- The eigenstates of the mass matrix are Majorana fermions (neutrinos \sim anti-neutrinos)

→ Smoking gun signature for Majorana neutrinos: neutrinoless double beta decay

- Eigenvalues of the mass matrix and the see-saw mechanism

$$m_{1,2} = \frac{m_R}{2} \pm \frac{1}{2} \sqrt{m_R^2 + 4m_D^2} \quad m_R \gg m_D \Rightarrow \begin{aligned} |m_1| &\approx m_R \\ |m_2| &\approx \frac{m_D^2}{m_R} \end{aligned}$$

* Searches at various experiments:
NEMO-3;
NEXT-100; KamLAND-Zen;
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Majorana Neutrinos and See-Saw Mechanism

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→ The mass of the light neutrino gets suppressed by a factor $m_D/m_R \ll 1$

Majorana Neutrinos and See-Saw Mechanism

Can we accommodate such scenario within asymptotically safe gravity?

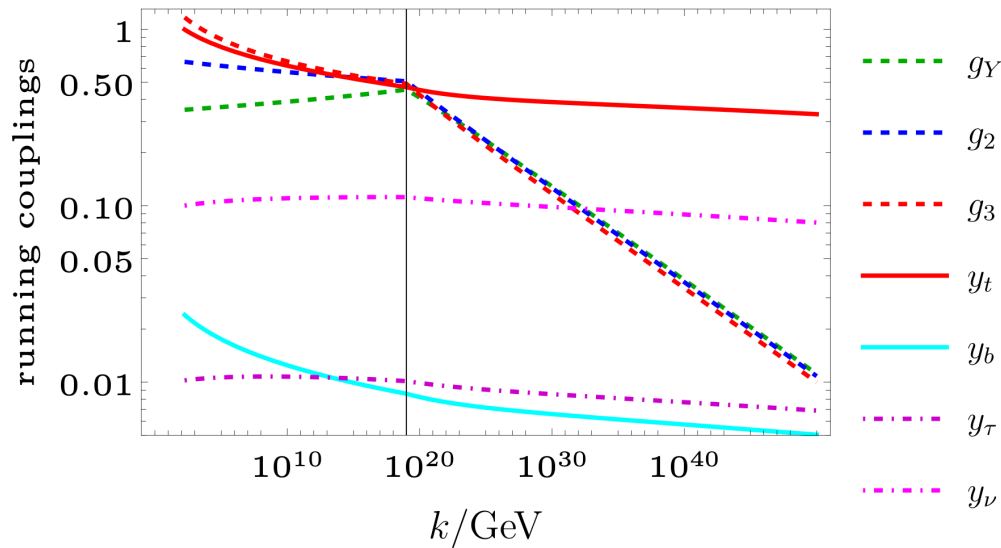
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- Yes! We can construct explicit RG trajectories that are UV complete for all couplings

➡ The Majorana mass is relevant at a fixed point with $m_R^* = 0$

Thus, we have enough freedom to accommodate non-vanishing Majorana masses in the infrared



Majorana Neutrinos and See-Saw Mechanism

Asymptotically safe gravity generates an upper bound on the see-saw scale

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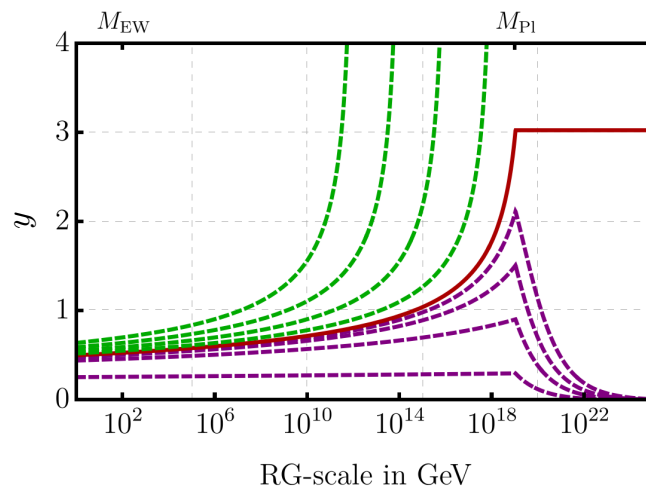
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Parametrised quantum gravity contribution

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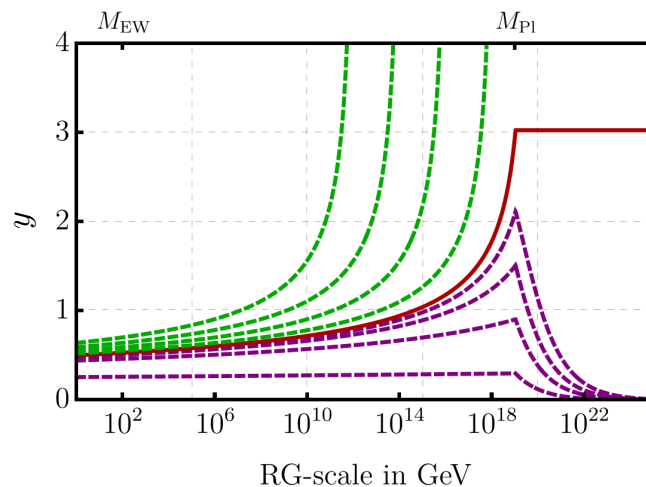
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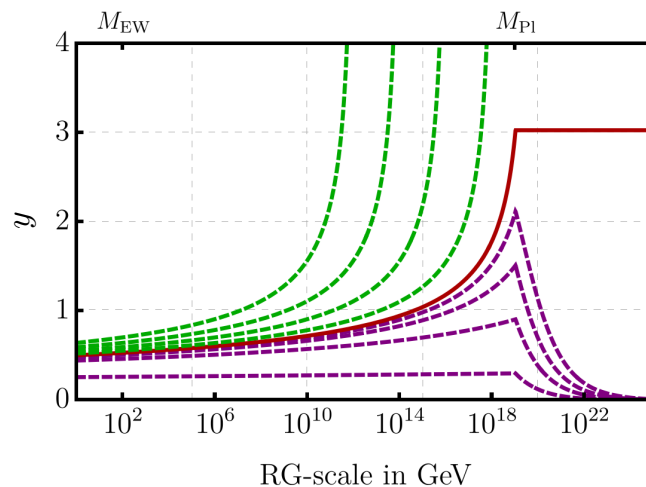
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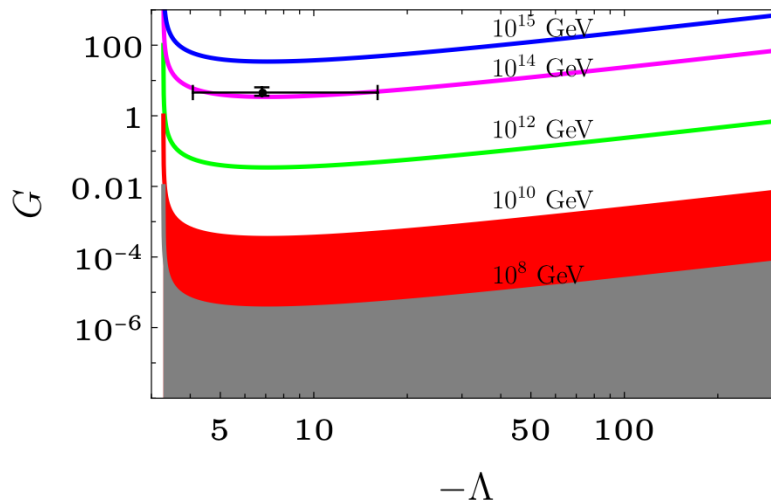
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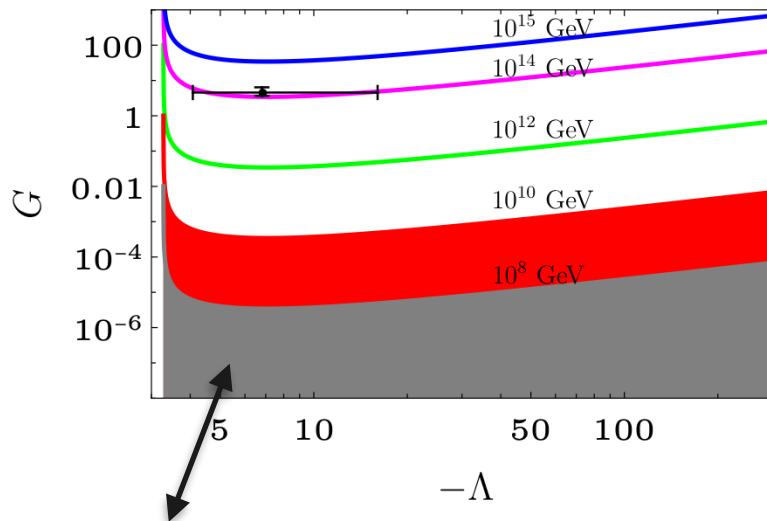
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Excluded if we impose from Davidson-Ibarra bound
(thermal leptogenesis)

Pseudo-Dirac Massive Neutrinos

- A different perspective on neutrinos with Majorana masses

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They behave almost like a Dirac neutrino, but not exactly

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- ➔ Almost maximal mixing between left-handed and right-handed neutrinos (*Dirac neutrinos features maximal mixing*)
- ➔ Oscillation between both chiral components, even with a single generation (*Dirac do not oscillate in the single generation case (degenerate states)*)

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* Experimental searches at JUNO and DARWIN

Franklin, Perez-Gonzalez and Turner, 2023
de Gouvêa, et. al., 2022

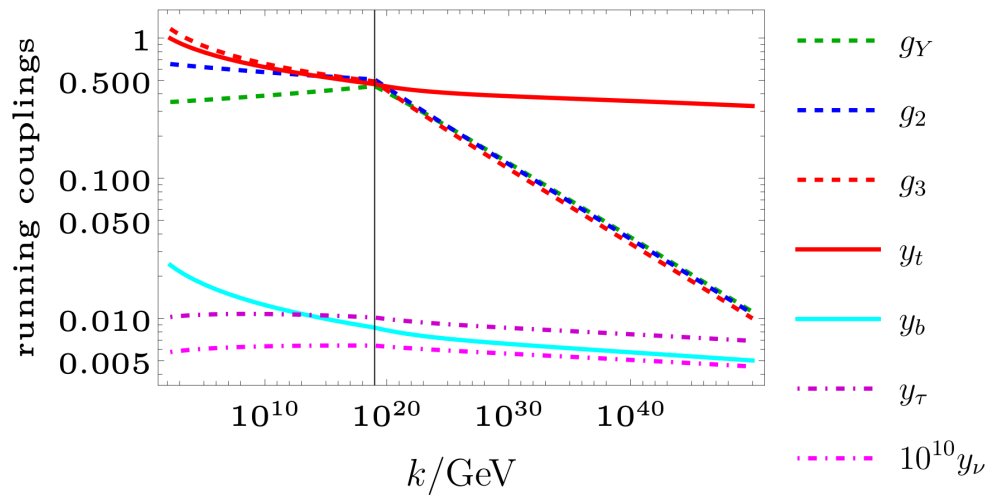
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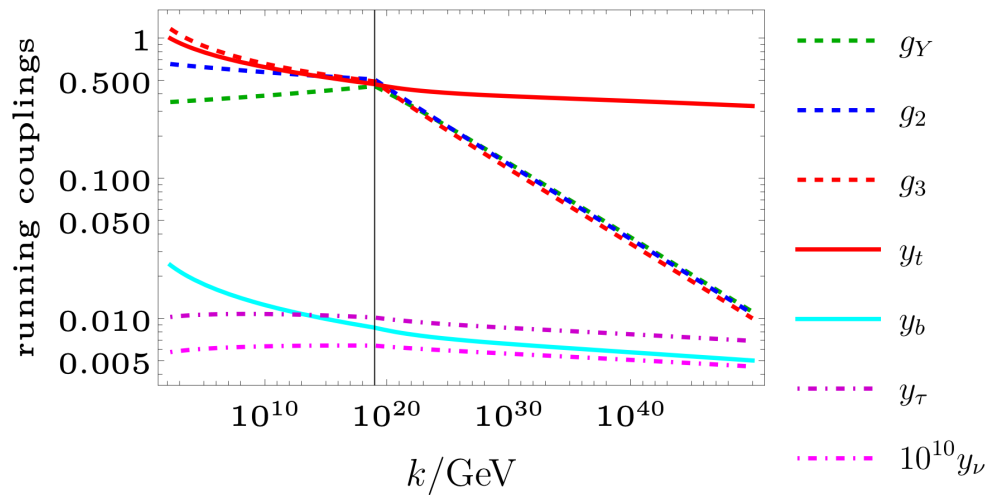
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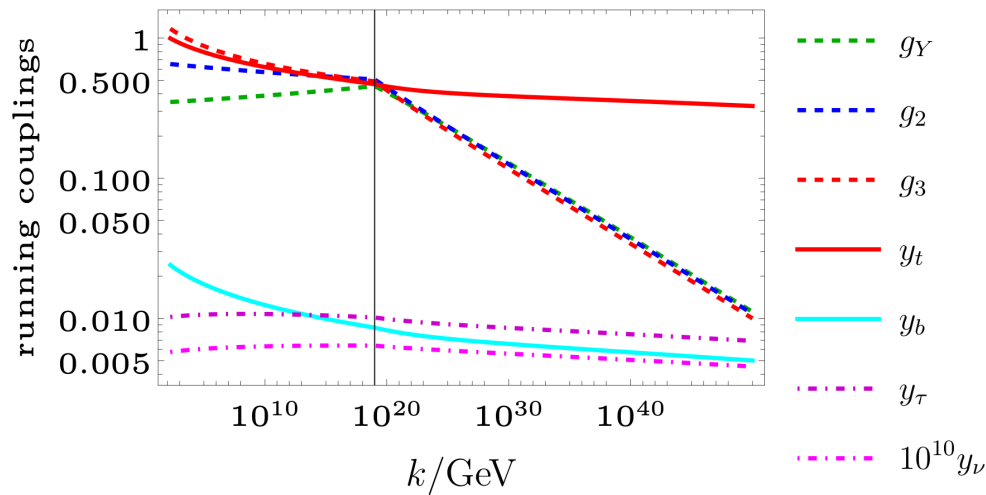
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➔ This scenario does not allow us to extract new theoretical bounds from asymptotically safe gravity



Summary

The interplay between massive neutrinos can give valuable information about the landscape of asymptotically safe gravity



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Pseudo-Dirac massive neutrinos

- There is enough room to tune the free parameters in harmony with asymptotic safety
- No upper or lower bound from this scenario

Thank you for your attention!